# Quantum (Second Harmonic) Efficiency and Conversion Coefficient for a Frequency Doubled He- Ne Laser. 

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#### Abstract

It is the aim of this project to study the ability of producing second harmonic generation for a ( 10 mW ) $\mathrm{He}-\mathrm{Ne}$ laser and to test some optical properties for KDP crystal and to compare the results with those obtained in testing the same crystal with using a ( 18 mW ) Laser Tube. It seems that the performance of visible radiation is more accurate rather than near IR radiation.


Keywords: Non- Liner Conversion Coefficient, Power Efficiency ( $\mathrm{P}_{2} / \mathrm{P}_{1}$ ), Second Harmonic Generation (SHG).

## 1. Introduction

Second harmonic generation is a non- linear response in those media which exhibits birefringence operation whose intensity reaches almost $\left(10^{8} \mathrm{~V} / \mathrm{m}\right)$. In non-linear media, pyroelectric materials show an electric polarization when the Temperature of a crystal is changed ${ }^{[1]}$. Thus a second Harmonic frequency of the light is changed from $(\omega)$ to $(2 \omega)$, if the crystal has an induced electric polarization owning a component oscillating at $(2 \omega)^{[2]}$. The relation between the polarization vector and the field strength in those media is:

$$
\begin{equation*}
\vec{P}=\epsilon_{\circ} x_{1} E+\epsilon_{\circ} x_{2} E_{2}^{2}+\epsilon_{\circ} \quad x_{3} E^{3}+ \tag{1}
\end{equation*}
$$

Where $\epsilon_{\circ}$ is the permittivity constant at space, $(\chi)_{s}$ are the non- linear susceptibility of the material through successive harmonic generation in the process ${ }^{[3]}$. Maxwell Equation for the non-linear media has a relative permeability $\in=k \in \in_{0}$, where k is the dielectric strength of the material.
Thus:

$$
\begin{equation*}
\nabla^{2} E=\mu \circ \in \frac{\partial^{2} E}{\partial t^{2}}=\mu_{\circ} \frac{\partial^{2}}{\partial t^{2}}(\in \circ \mathrm{E}+\mathrm{P}) \tag{2}
\end{equation*}
$$

Here P is again the polarization vector. For first harmonic wave equation all polarization components oscillating at ( $2 \omega$ ) must be induced to equation (1), as such ${ }^{[4]}$ :

$$
\nabla^{2} E_{2}=\mu \circ \in \circ \frac{\partial^{2}}{\partial t^{2}}\left(E \circ+x_{1} E_{2}+x_{2} E_{2}^{2}+\right.
$$

$\qquad$
Where this recognizes the refractive index at second harmonic generation as, $\mathrm{n}_{1}$ and $\mathrm{n}_{2}$, where $\mathrm{n}_{2}$ is related to the susceptibility as:

$$
\begin{equation*}
n_{2}^{2}=1+x_{1} \tag{4}
\end{equation*}
$$

Putting this in equation (3), we have:

$$
\begin{equation*}
\nabla^{2} E_{2}=\mu_{\circ} \in_{\circ} \frac{\partial^{2} E}{\partial t^{2}}\left(n_{2}^{2}+x_{2} E_{1}^{2}\right) \ldots \ldots \ldots \ldots \ldots \ldots \ldots \tag{5}
\end{equation*}
$$

The detail of mathematical formalism take long and the second harmonic power of the light beam [if Laser] is related to the fundamental beam power as:

$$
\begin{equation*}
P_{2}=\sqrt{\frac{\mu_{\circ}}{\epsilon_{\circ}}} \frac{\omega^{2} x_{2}^{2} l^{2}}{8 n_{1}^{2} n_{2}} P_{1}^{2} \tag{6}
\end{equation*}
$$

Remembering that in classical EM theory: $P=\frac{E^{2}}{\nabla t}$ Where $E^{2}$ is the spot intensity of the light.
Here $n_{1}$ is called the fundamental refractive index and $n_{2}$ is second harmonic refractive index, occasionally, they are called [ordinary and extra ordinary refractive indices]. Making a few approximations to insert another parameter which the constant of the harmonic generation, related to the non-linear conversion coefficient $\left(\mathrm{d}_{\text {ooe }}\right)$ of the crystal as:

$$
\begin{equation*}
K=\frac{128 \pi^{2} \omega_{1}^{2}\left(d_{\left.00 e^{+}+\right)^{2}}^{n_{1} c^{3}} \sin ^{2} \theta_{m},{ }^{2}\right.}{} \tag{7}
\end{equation*}
$$

Here $\theta_{\mathrm{m}}$ is called the birefringence angle of the crystal, obtained from Brags diffraction law analysis for the crystal where value is ${ }^{[6]}$ :

$$
\begin{equation*}
\sin \vartheta_{m}=\frac{\left(n_{2}^{2 \omega_{1}}\right)^{2}\left[\left(n_{1}^{2 \omega_{1}}\right)^{2}-\left(n_{1}^{\omega_{1}}\right)^{2}\right]}{\left(n_{1}^{\omega_{1}}\right)^{2}\left[\left(n_{1}^{2 \omega_{1}}\right)^{2}-\left(n_{2}^{2 \omega_{1}}\right)^{2}\right]} \tag{8}
\end{equation*}
$$

Combining equations (8) and (6) we get ${ }^{[7]}$ :

$$
\begin{equation*}
P_{2}=K P_{1}^{2} \frac{l . l_{a}}{\omega_{\circ}^{2}} \tag{9}
\end{equation*}
$$

Here 1 is the crystal length, $\omega \circ$ is the fundamental beam spot size and $l_{a}$ is the laser aperture beam, related to the phase matching angle $\left(\theta_{m}\right)$ as ${ }^{[8]}$ :

$$
\begin{equation*}
l_{a}=\sqrt{\pi} \frac{\omega \circ}{p} . \tag{10}
\end{equation*}
$$

Putting this in equation (9), we obtain:

$$
\begin{equation*}
P_{2}=K P_{1}^{2} \sqrt{\pi} \frac{l}{\omega \circ \rho} \ldots \tag{11}
\end{equation*}
$$

This will be the key equation for our calculations, since it combines both fundamental beam power [ $P_{1}$ ] and the second harmonic power $\left[P_{2}\right]$ and $(\rho)$ is related to $\theta_{m}$ as:

$$
\begin{equation*}
\tan p=\frac{1}{2}\left(n_{1}^{\omega_{1}}\right)^{2}\left[\frac{1}{\left(n_{2}^{2 \omega}\right)^{2}}-\frac{1}{\left(n_{1}^{\left(2 \omega_{1}\right)}\right)^{2}}\right] \sin ^{2} \theta_{m} \tag{12}
\end{equation*}
$$

## 2- Results and Calculations

In the process, a birefringence crystal which is shown in figure (1) used having the following properties [9]:
1- Geometric dimensions: $\mathrm{h}=0.5 \mathrm{~cm}, \mathrm{l}=2 \mathrm{~cm}, \omega=7 \mathrm{~cm}$
2- Refractive index: $\mathrm{n}_{-}^{\circ}=1.5072$,n_(e )=1.492
3- Optical properties:
a. $\quad$ d_00e $=1.316 \times \llbracket 10 \rrbracket \wedge(-11) \mathrm{m} / \mathrm{V} \cdots$. .
b. phase matching angle $(\theta \mathrm{m})=(56.1)^{\circ} \cdots \cdots \cdots$.
c. birefringence angle [Rad.] $=0.0281$

From these and for the ordinary light beam for a Typical He-Ne Laser ( $\lambda=632.8 \mathrm{~nm}$ ), we could calculate the non-linear conversion coefficient $K$ eqn. (9) to be $2.632 \times \llbracket 10 \rrbracket \wedge(-11)(\mathrm{m} / \mathrm{V})$. Having doing these, it is an easy matter to plot both ( $\mathrm{P} \_2$ ) and Second Harmonic efficiency $\rho_{-}$SH as the function of the fundamental beam power for available output (He-Ne) Laser powers ranging from (1-30)mW. Figure (1) shows the graph between the Second Harmonic power (P2) and the fundamental beam power P. Figure (2) shows variation of Second Harmonic efficiencies ( $\mathrm{P} \_2 / \mathrm{P} \_1$ ) as the function of $\lambda$.

## 3-Conclusions

Throughout the whole work, we observed that the new version of calculating the conversion efficiency and second Harmonic power have been prompted potentially. The version used was based on the idea of using low power He-Ne Laser to enhance mechanical and thermal effects. The figures show that:
1- The second Harmonic power is more stable at fundamental beam power ranging from (1-10) m Watts and the Fluctuation of raising up is starting nearly (20)mV.
2- 2- The Second Harmonic power is much efficient at near uv(360-4000nm beam wavelength reaching its mid value at (650)nm which is almost close to the employed wave length of the laser radiation.
3- The present work have differentiated the Second Harmonic power as the function of refractive index from that of Storelu, by fixing the value of the non-Linear Conversion efficiency $[\mathrm{K}]$ and observing the changes in geometrical parameters of the KDP crystal with in the optical properties allowed.

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Figure1. A birefringence crystal dimensions.


Figure2. The graph between the second Harmonic PowerP $P_{2}$ and the fundamental beam $\mathrm{P}_{1}$.


Figure 3. The variation of second Harmonic Efficiencies $\left(\mathrm{P}_{2} / \mathrm{P}_{1}\right)$ as a function of $\lambda$.

